



A1 Collaboration

presented by

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Mainz, Bonn, Darmstadt, Glasgow, Ljubljana, Zagreb

Measurement of the Electric Form Factor of the Neutron at MAMI

- Introduction
- Measurement of $G_{E,n}$ in $D(\vec{e}, e' \vec{n})p$
- The A1 Experiment
- Data Analysis
- Summary

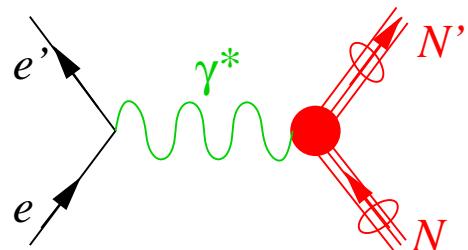
Spin 2002
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Upton, NY, September 9-14, 2002

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Nucleon Form Factors

Form factors parametrize finite size of a particle

- Electromagnetic current:



$$\bar{N}_{\text{point}} \gamma_\mu N_{\text{point}} \rightarrow \bar{N} \Gamma_\mu N$$

$$= \bar{N} [\gamma_\mu F_1(Q^2) + \frac{i\sigma_{\mu\nu}q^\nu}{2M_N} \kappa F_2(Q^2)] N$$

$F_1(Q^2)$: Dirac form factor

$F_2(Q^2)$: Pauli form factor

- Sachs form factors:

$$G_E(Q^2) = F_1(Q^2) - \tau \kappa F_2(Q^2)$$

$$G_M(Q^2) = F_1(Q^2) + \kappa F_2(Q^2)$$

with $\tau = Q^2/4M_N^2$.

$G_{E,n}$ Measurements

Difficulties in measurement of $G_{E,n}$:

- No free neutron target available; use of D, ${}^3\text{He}$ in general model dependent
- Rosenbluth cross section:

$$\frac{d\sigma}{d\Omega} = \left(\frac{d\sigma}{d\Omega} \right)_M \cdot \frac{E'_e}{E_e} \cdot \left(\frac{G_E^2 + \tau G_M^2}{1 + \tau} + 2\tau G_M^2 \tan^2 \frac{\vartheta_e}{2} \right)$$

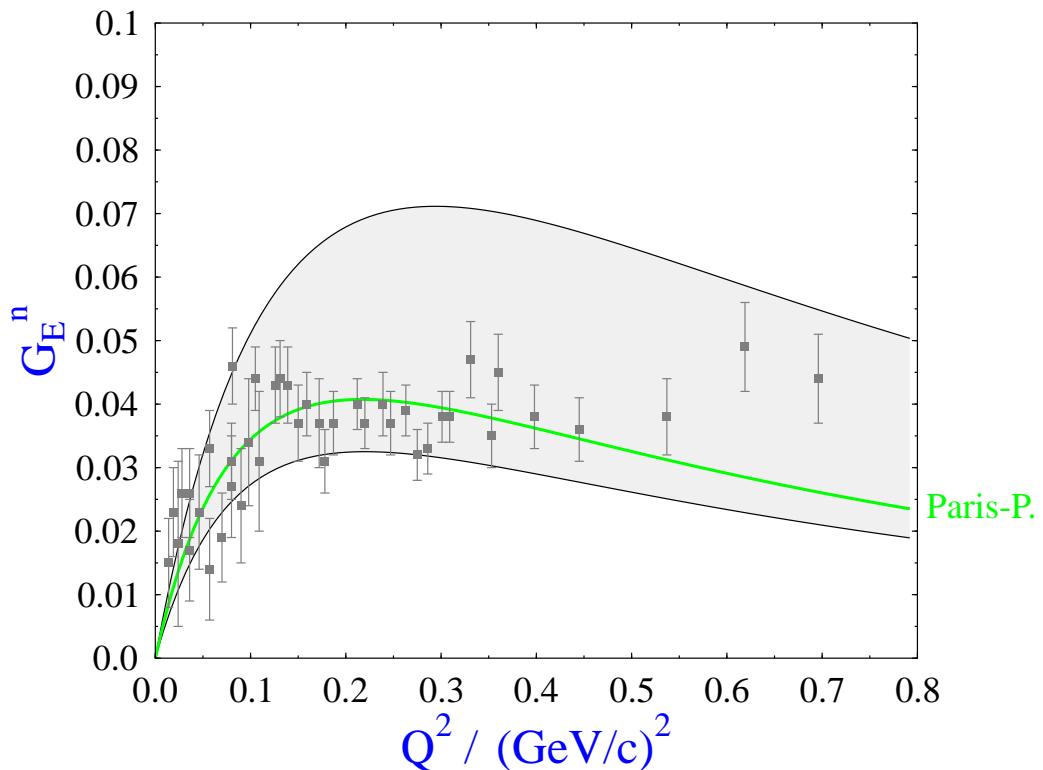
$G_{E,n}^2 \ll \tau G_{M,n}^2 \Rightarrow$ Rosenbluth separation difficult

However, $G_{E,n} \neq 0$ at finite Q^2 , as proved in two kinds of experiments:

- $\langle r_{ch,n}^2 \rangle = -0.1148(23) \text{ fm}^2 / -0.1243(28) \text{ fm}^2$
(S. Kopecky et al., PRC **56** (1997), 2229)
- $G_{E,n}$ from $D(e, e')d$
(S. Platchkov et al., Nucl. Phys. **A510** (1990), 740)

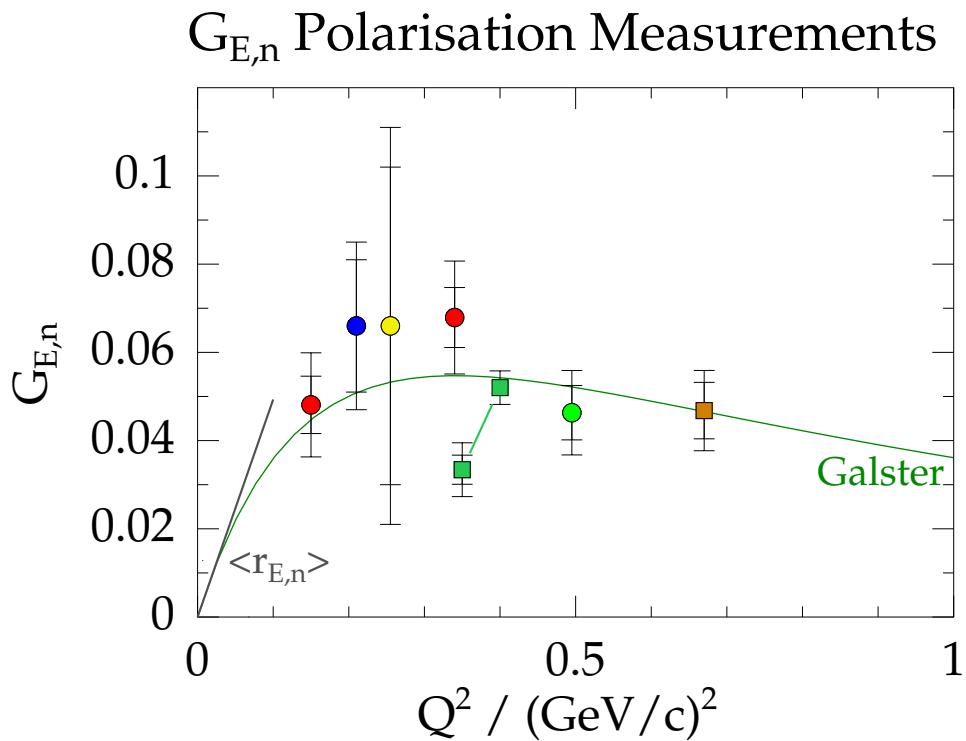
$G_{E,n}$ Measurements

Platchkov et al.:



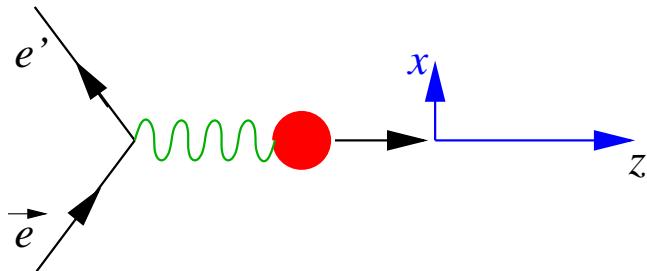
- $D(e, e')d$ model dependent!
- Model-independent results achieved in double polarization experiments

$G_{E,n}$ Data



- $D(\vec{e}, e' \vec{n})p$ at MIT-Bates, 1994
- $D(\vec{e}, e' \vec{n})p$ at MAMI (A3), 1999
- $\bar{D}(\vec{e}, e' n)p$ at NIKHEF, 1999
- ${}^3\vec{\text{He}}(\vec{e}, e' n)pp$ at MAMI (A3), 1999
- $\bar{D}(\vec{e}, e' n)p$ at JLab, 2001
- ${}^3\vec{\text{He}}(\vec{e}, e' n)pp$ at MAMI (A1), 2001

Measurement of $\sigma_{E,n}$ in $D(\vec{e}, e'\vec{n})p$

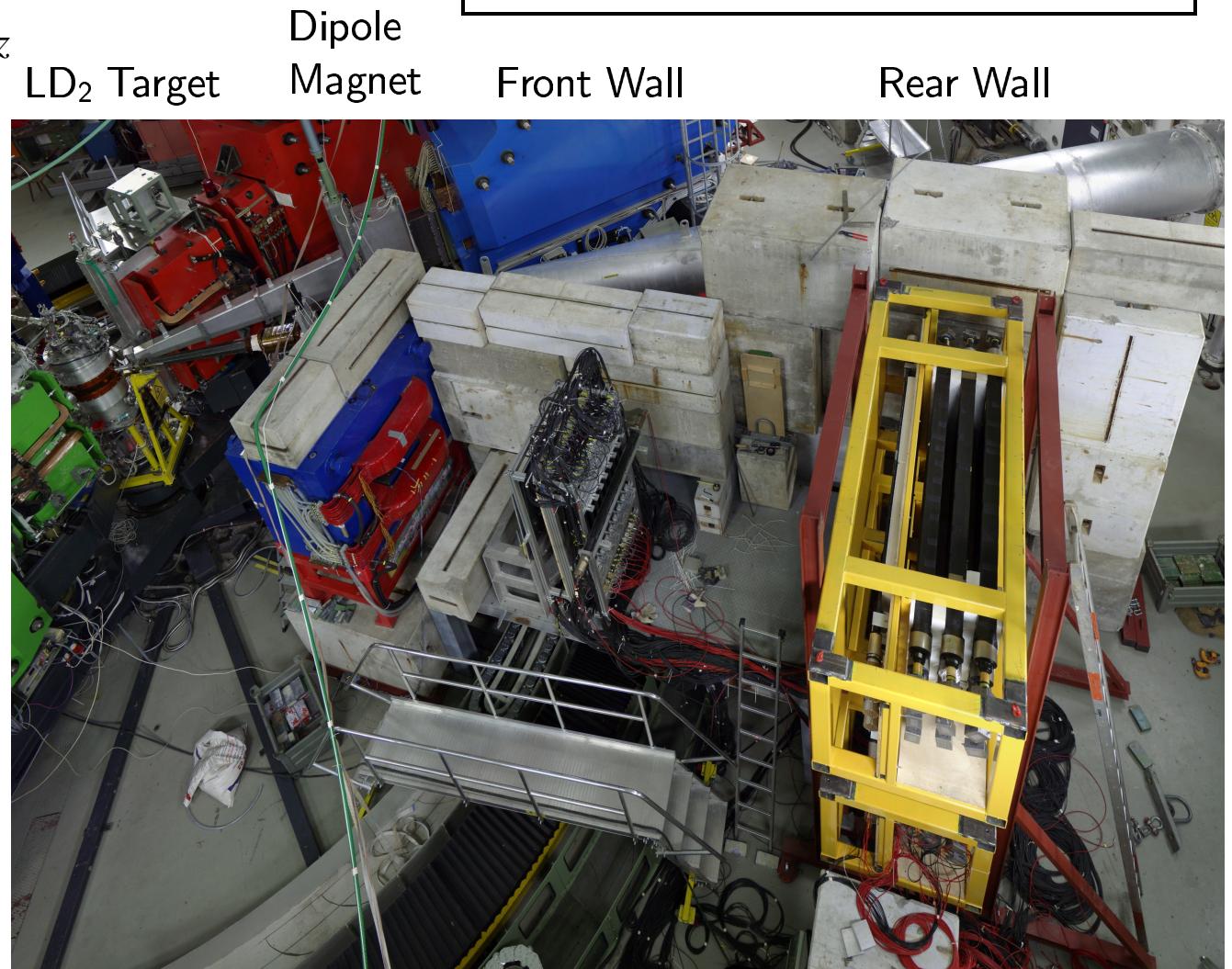


$$\text{Asymmetry } A = P_e \mathcal{A}_{\text{eff}} P_t \sin \Phi_n$$

$$\begin{aligned}\mathcal{P}_x &= -hP_e \frac{aG_E G_M}{G_E^2 + bG_M^2} \\ \mathcal{P}_y &= 0 \\ \mathcal{P}_z &= hP_e \frac{cG_M^2}{G_E^2 + bG_M^2}\end{aligned}$$

Arnold, Carlson & Gross,
PR C 23 (1981), 363

Method:
T. N. Taddeucci et al.,
NIM A241 (1985), 448



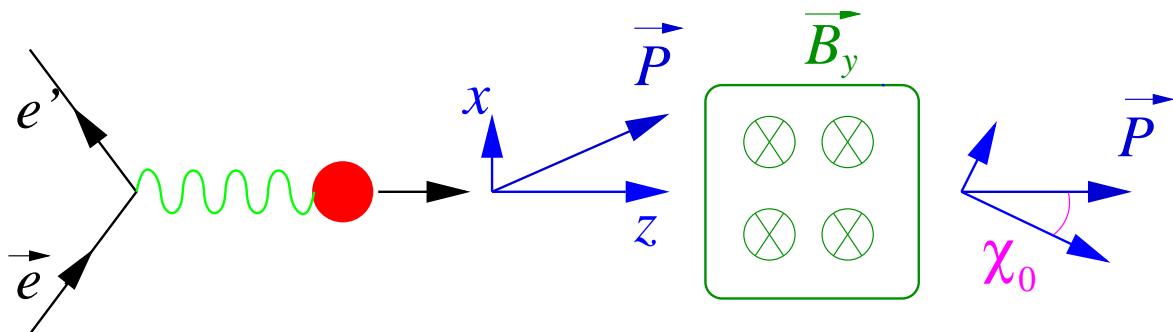
Measurement of $G_{E,n}$ in $D(\vec{e}, e' \vec{n}) p$

Spin Precession Method

Mixing of two spin components by precession of neutron spin in magnetic field,

$$\mathcal{P}_t(\chi) = \mathcal{P}_x \cos(\chi) + \mathcal{P}_z \sin(\chi) =: \mathcal{P}_0 \sin(\chi - \chi_0)$$

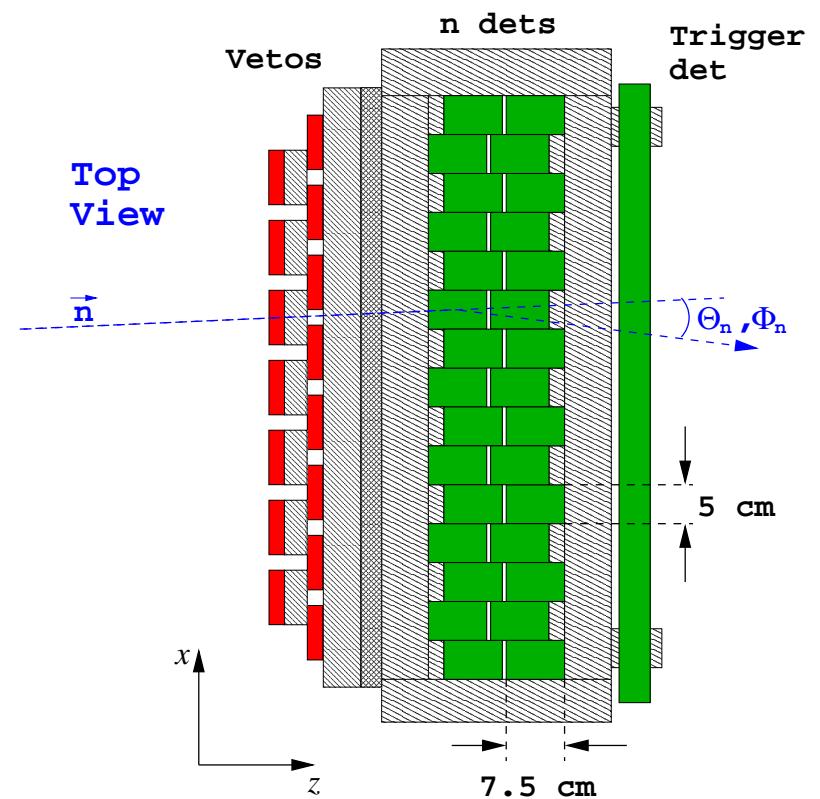
(M. Ostrick et al., PRL **83** (1999) 276)



$$\tan \chi_0 = \frac{\mathcal{P}_x}{\mathcal{P}_z} = \frac{A_x}{A_z} = \frac{a \mathcal{A}_{\text{eff}} P_e}{c \mathcal{A}_{\text{eff}} P_e} \cdot \frac{G_{E,n}}{G_{M,n}} .$$

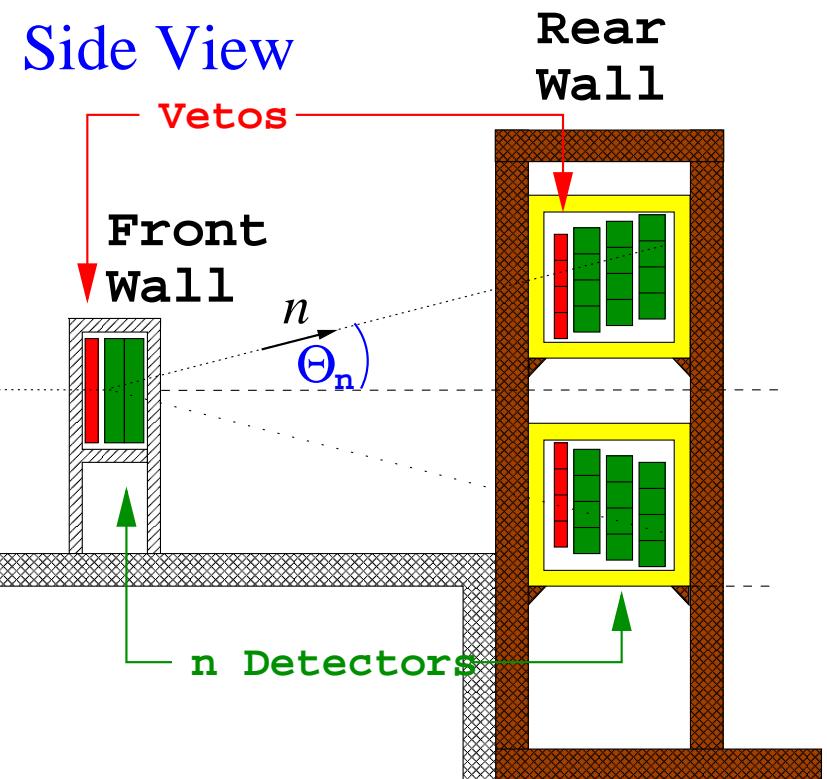
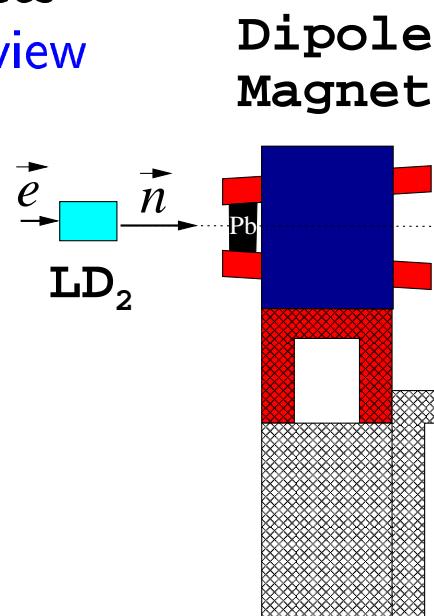
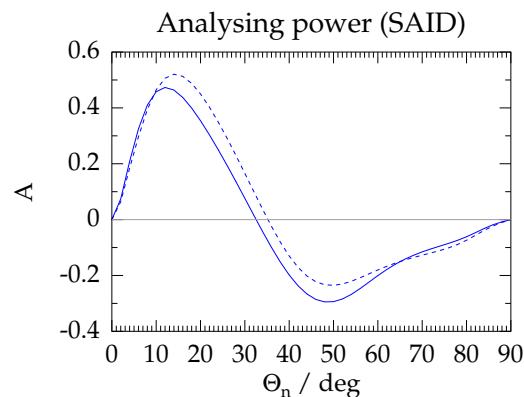
Neutron Polarimeter, Front Wall

- Overall size: $80 \times 80 \text{ cm}^2$
- 5 cm scintillator width
at 6 m target distance
 \Rightarrow Angular resolution 0.48° FWHM
- Charged particle vetoing
- Single rates $\simeq 1 \text{ MHz}$
- 15% neutron detection efficiency



Neutron Polarimeter, Rear Wall

24 scintillator bars, $180 \times 20 \times 10 \text{ cm}^3$
 ⇒ high background rates
 ⇒ avoid direct target view



Data Taking

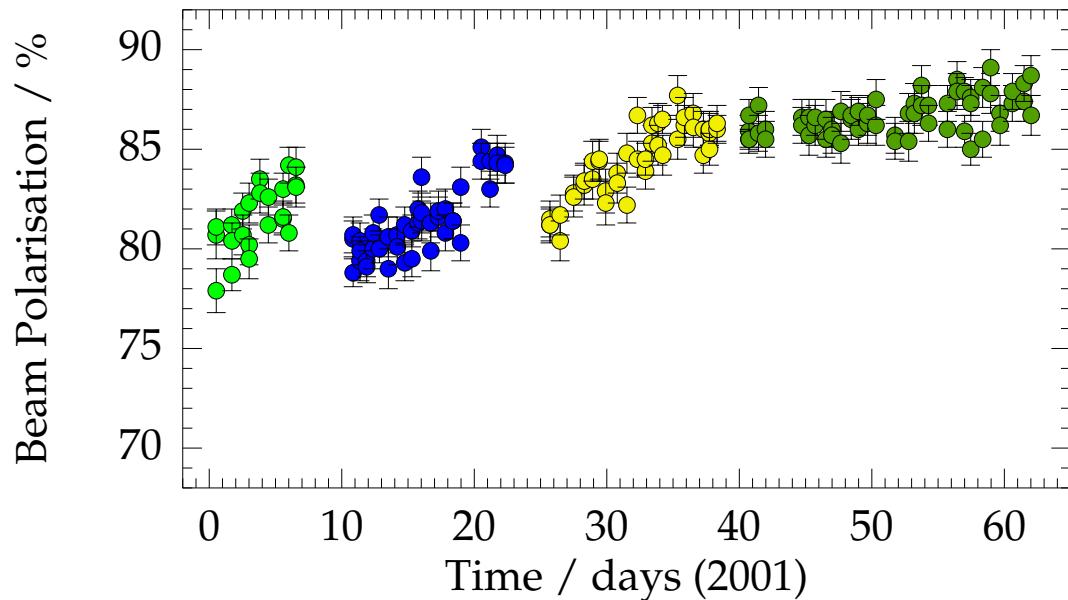
- Three central momentum transfers:

$Q^2 / (\text{GeV}/c)^2$	0.6	0.8	0.3	
E_e / MeV	855	883	660	
E'_e / MeV	536	454	498	Spectrometer A
ϑ_e	70°	90°	57°	
ϑ_n	37°	27°	47°	Neutron polarimeter
T_n / MeV	320	427	160	

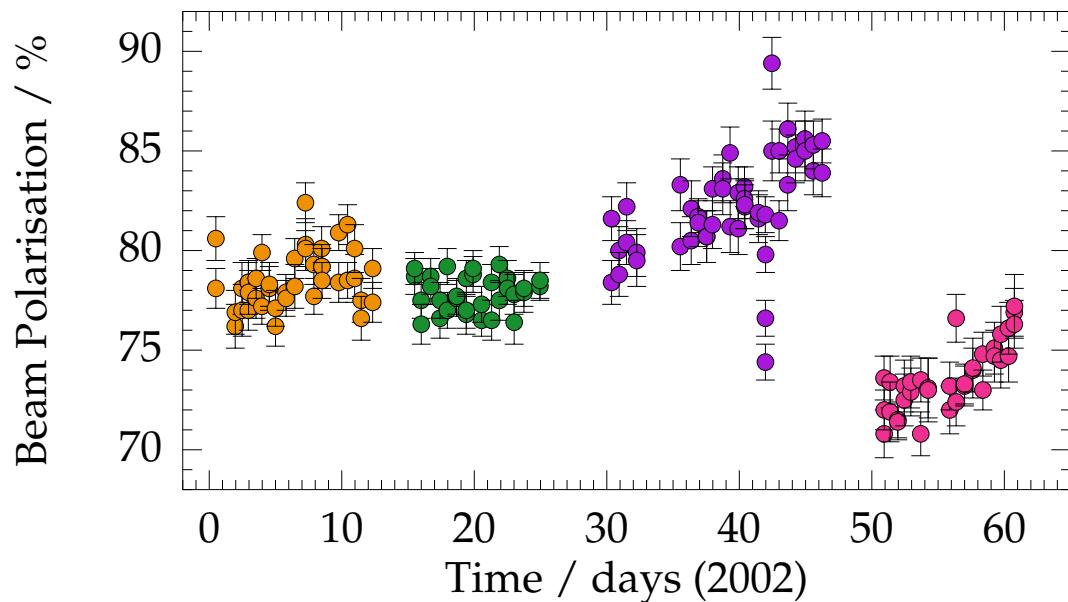
- LD₂ target (length: 5 cm)
 - Beam current: $\sim 10 - 15 \mu\text{A}$
 - 5 cm lead shielding
 - LH₂ data (calibration, proton misidentification)
 - Beam polarisation monitored with Møller polarimeter*, $P_e \simeq 80\%$
- *P. Bartsch, PhD thesis, Mainz 2001

Electron Beam Polarisation

Moller Polarisation Measurements



Moller Polarisation Measurements



Data Analysis

Outline:

- Adjustment of all scintillators (\Rightarrow calibration)
- Reconstruction of reaction kinematics
- Neutron identification, background suppression
- Calculation of beam-helicity asymmetries
- Correction of asymmetry shifts due to
 - Variations of electron beam polarisation
 - p-n conversion in lead shielding
 - Nuclear binding effects

Data Analysis

Reaction Kinematics

Complete reconstruction of quasi-elastic $e - n$ scattering requires 5 observables:

Spectrometer A		n polarimeter	
ϑ_e, φ_e E'_e	$< 2 \text{ mrad}$ 10^{-4}	ϑ_n, φ_n ToF_n	8 mrad $< 1 \text{ ns}$

In Spectrometer A:

- Preselection of quasielastic events via $e - n$ coincidence time
- Further background reduction, e.g. pion suppression with Cerenkov detector

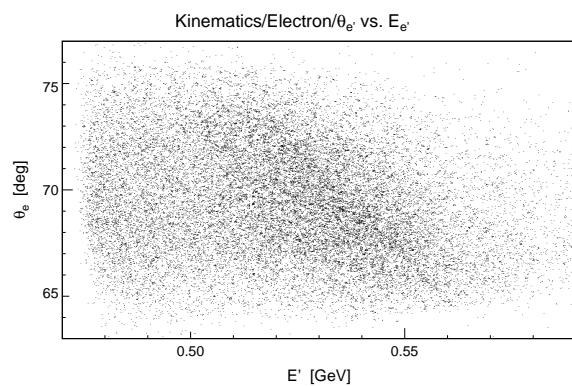
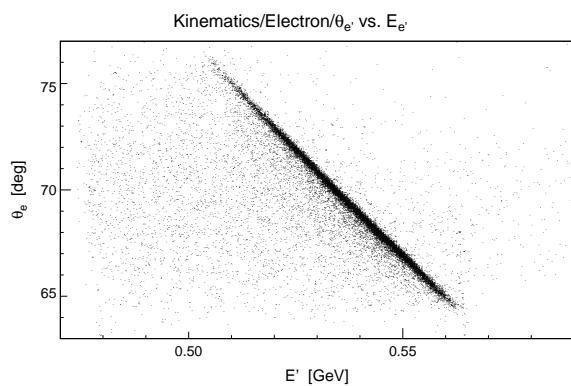
Data Analysis

Effect of cuts on Spectrometer A

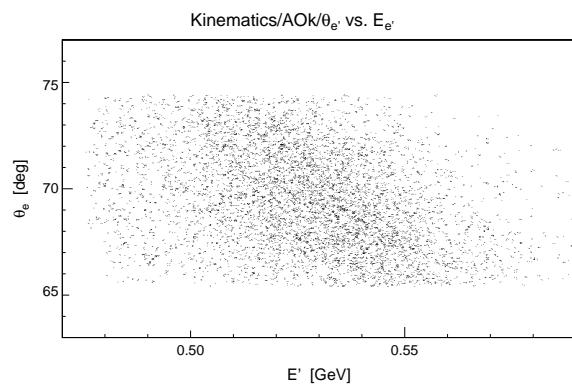
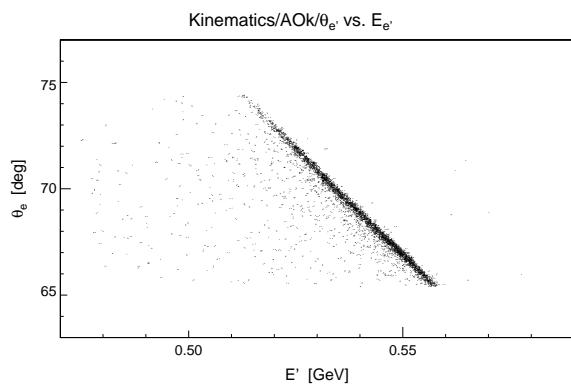
Proton (LH_2) data

Neutron (LD_2) data

Elastic line in focal plane:



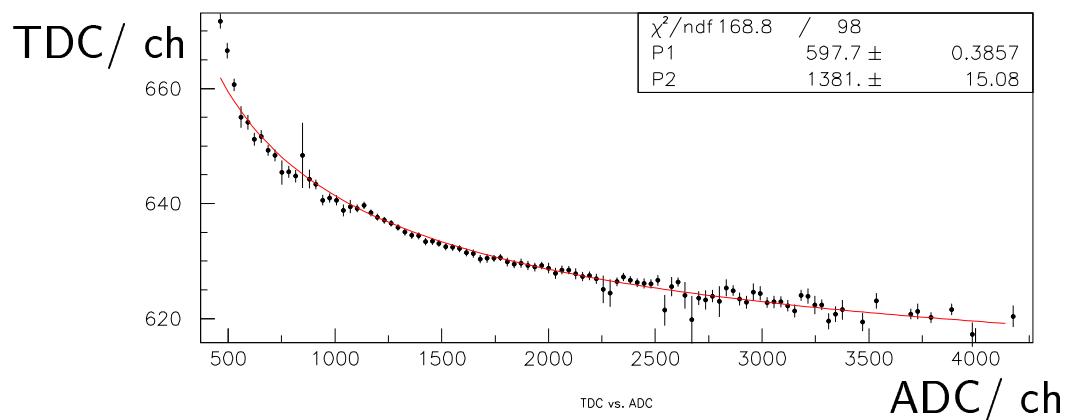
After cuts on Spectrometer A:



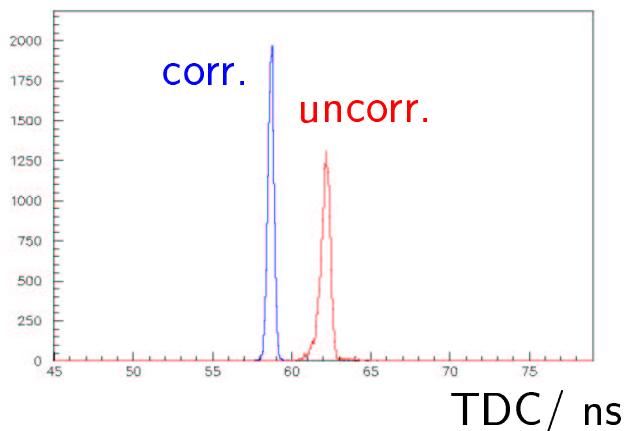
Calibration

Time of Flight

Walk correction (using electrons):



Result:



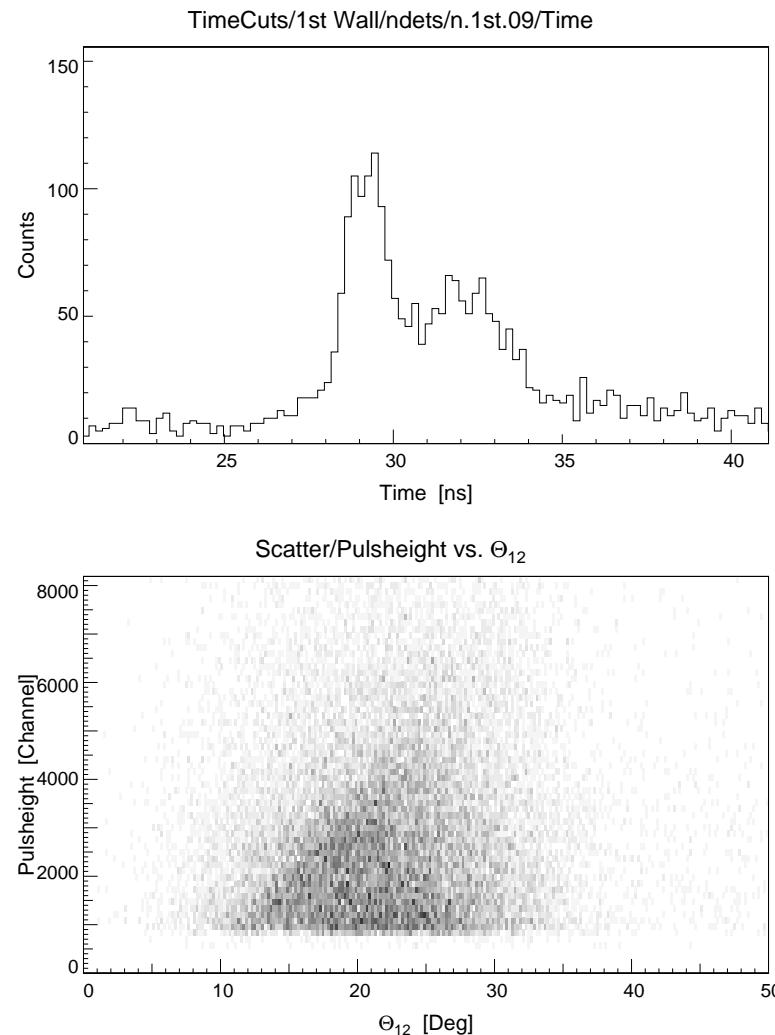
Resolution in neutron detectors: $\sigma_t = 300 \text{ ps}$
 \Rightarrow Spatial resolution: $\sigma_x = 4.5 \text{ cm}$
 \Rightarrow Angular resolution: $\sigma_\vartheta = 7.5 \text{ mrad}$

Data Analysis

Neutron Identification

Suppress proton background via

- Charged particle vetoing
- ToF: protons lose kinetic energy in lead shield
- Correlation between energy deposition of neutrons in 1st wall and scattering angle Θ_n

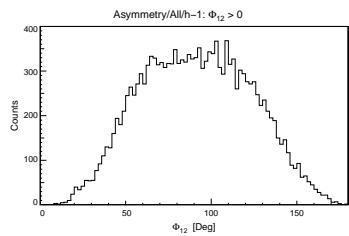


Data Analysis

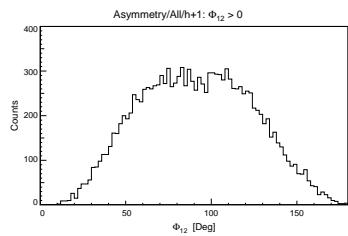
Asymmetries

Φ_n distributions:

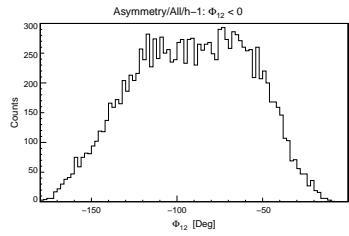
$h = -1, \Phi_n > 0$



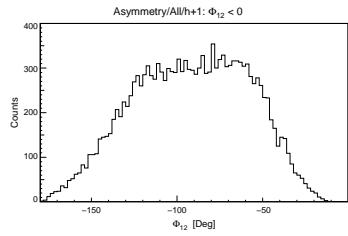
$h = +1, \Phi_n > 0$



$h = -1, \Phi_n < 0$

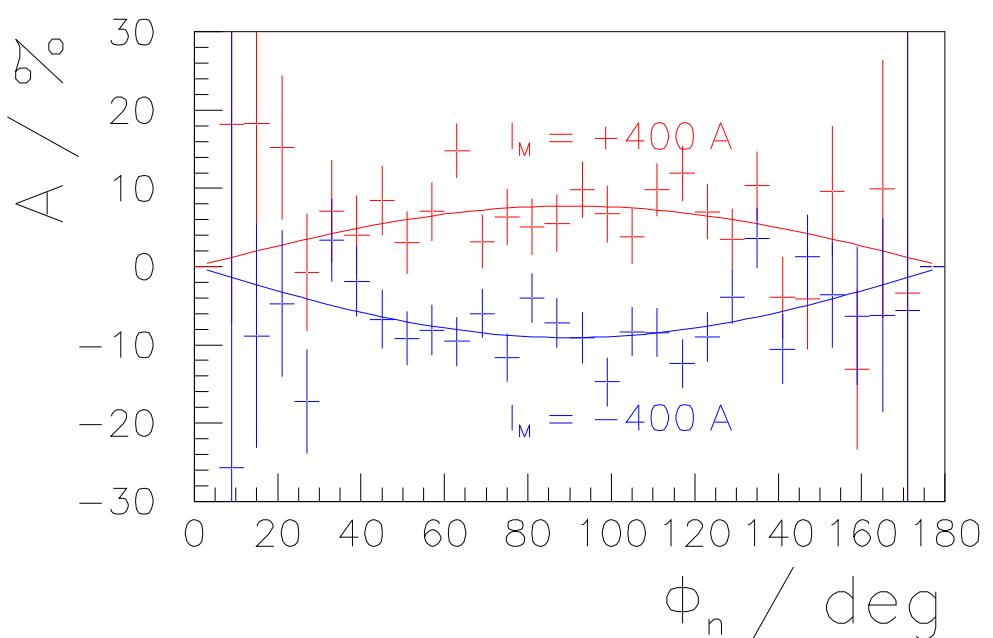


$h = +1, \Phi_n < 0$



$$A(\Phi_n) = \frac{\sqrt{N^+(\Phi_n)N^-(\Phi_n+\pi)} - \sqrt{N^+(\Phi_n+\pi)N^-(\Phi_n)}}{\sqrt{N^+(\Phi_n)N^-(\Phi_n+\pi)} + \sqrt{N^+(\Phi_n+\pi)N^-(\Phi_n)}}$$

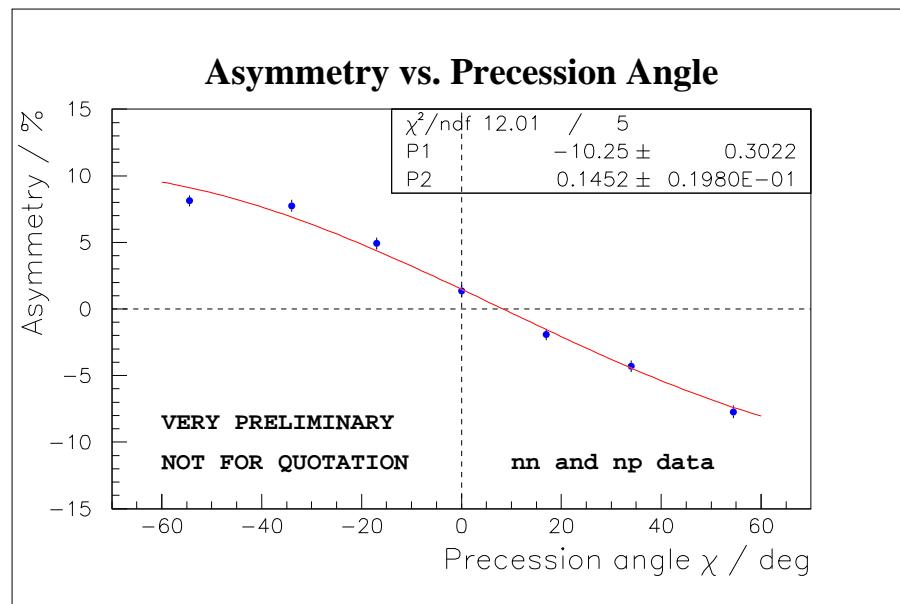
Asymmetry



Michael Seimetz

Asymmetries

Plot asymmetries as function of spin precession angle χ



Summary

- $G_{E,n}$ is a key observable in hadron physics
- Double polarization experiments allow precise, model-independent $G_{E,n}$ measurements
- New $D(\vec{e}, e' \vec{n})p$ experiment in the A1 collaboration at Mainz
- $Q^2 = 0.3, 0.6$ and $0.8 (\text{GeV}/c)^2$
- Data taking completed in August 2002
- Data analysis in progress, first promising results available