

Determination of the Gamow-Teller Quenching Factor via the $^{90}\text{Zr}(n, p)$ Reaction at 293 MeV

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Gamow-Teller (GT) resonances have been extensively studied since its discovery in 1975. The GT transition is characterized as spin-flip ($\Delta S = 1$), isospin-flip ($\Delta T = 1$) and no transfer of orbital angular momentum ($\Delta L = 0$). There exists a model-independent sum rule, $S_{\beta^-} - S_{\beta^+} = 3(N - Z)$, where S_{β^-} and S_{β^+} are the GT strength of β^- and β^+ types, respectively [1]. Surprisingly, however, only a half of the GT sum rule value was identified from the (p, n) measurement on targets throughout the periodic table [2]. This problem, so-called the quenching of the GT strengths, has been one of the most interesting phenomena in nuclear physics because it is related to non-nucleonic (Λ -isobar) degrees of freedom in nuclei; the quenching factor sets a strong constraint on the Landau-Migdal parameters, g'_{NN} and g'_{NA} , in the $\pi + p + g'$ model [3].

Recently, Wakasa et al. have measured the angular distribution of the double differential cross sections for the $^{90}\text{Zr}(p, n)$ reaction at 295 MeV. By performing multipole decomposition (MD) analysis, the GT strengths of $S_{\beta^-} = 28.0 \pm 1.6$ has been obtained in the continuum up to 50 MeV excitation in ^{90}Nb [4]. Determination of the Δ -isobar contribution to the GT sum rule, however, requires precise (n, p) cross-section data at the same energy. For this purpose we have constructed in an (n, p) facility at RCNP and measured the double differential cross sections for the $^{90}\text{Zr}(n, p)^{90}\text{Y}$ reaction at 293 MeV.

We have obtained the differential cross sections up to 70 MeV excitation energy over an angular range of 0° - 12° with a statistical accuracy of $1.7\%/2 \text{ MeV}$ at 1° - 2° . Figure 1 shows the result of the MD analysis. After subtracting the contribution from the isovector spin monopole (IVSM) resonance observed at $E_x \approx 20 \text{ MeV}$ with a width of $\sim 10 \text{ MeV}$, we have obtained a total GT strength of $S_{\beta^+} = 3.0 \pm 0.3 \pm 0.8 \pm 0.5$ up to 31.4 MeV excitation where the errors are uncertainties of the MD analysis, the IVSM contribution, and the GT unit cross section.

By using the S_{β^-} value by Wakasa et al., the quenching factor Q , which is defined by $Q \equiv \frac{S_{\beta^-} - S_{\beta^+}}{3(N - Z)}$, has been

deduced to be $Q = 0.83 \pm 0.06$ in regard to Ikeda's sum rule value of $3(N - Z) = 30$. Therefore, the quenching of the GT strength due to the ΔN^{-1} admixture into the $1p1h$ GT state is significantly smaller than the quenching of $\sim 50\%$, observed in the previous studies [2] where the GT strengths in the continuum are not taken into account. Then the Landau-Migdal parameters, g'_{NA} and g'_{NN} , have been determined from the quenching factor. The deduction by Suzuki and Sakai [31] in Chew-Low model leads to $g'_{NN} \simeq 0.6$ and $0.16 < g'_{NA} < 0.35$ for $g'_{\Delta\Delta} = 0.6$

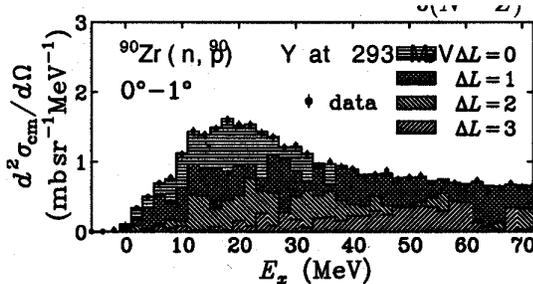


Figure 1: The result of MD analysis on the double differential cross sections for the $^{90}\text{Zr}(n, p)^{90}\text{Y}$ reaction at 293 MeV.

References

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